## Effects of dust feedback on evolutions of disks and planets

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## Viscous evolution of a gas disk

**Protoplanetary disk** 



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### Gas-dust friction vs. Viscous evolution

$$\begin{split} \Omega_{\rm gas} &= \Omega_{\rm K} \sqrt{1 - \eta} \qquad \eta = -\frac{1}{2} \left(\frac{h}{R}\right)^2 \frac{\partial \ln P}{\partial \ln R} \qquad \eta > 0 \text{ in a usual case} \\ \Omega_{\rm dust} &= \Omega_{\rm K} \end{split}$$



### **Basic equations**

Equations of motions of dust grains and disk gas For the dust grains Gas-dust friction



T<sub>stop</sub>: stopping time

For the disk gas  $\frac{\partial V_{\rm R}}{\partial t} + \left(\vec{V} \cdot \nabla\right) V_{\rm R} - \frac{V_{\phi}^2}{R} = -\frac{c_s^2}{\rho_{\rm g}} \frac{\partial \rho_{\rm g}}{\partial R} - \frac{GM_*}{R^2} + \frac{f_R}{\rho_{\rm g}}$   $\frac{\partial V_{\phi}}{\partial t} + \left(\vec{V} \cdot \nabla\right) V_{\phi} + \frac{V_{\rm R} V_{\phi}}{R} = -\frac{f_{\phi}}{\rho_{\rm g}}$ 

 $f_R$ ,  $f_{\phi}$ : viscous forces acting unit area in radial and azimuthal directions.

### Dust and gas velocities in steady state

• Gas velocity

Dipierro+17, Kanagawa+17

$$V_{\rm R}(R,z) = \frac{2St'}{St'^2 + 1} \frac{\rho_{\rm d}}{\rho_{\rm g} + \rho_{\rm d}} \eta V_K + \left(1 - \frac{1}{St'^2 + 1} \frac{\rho_{\rm d}}{\rho_{\rm g} + \rho_{\rm d}}\right) V_{\rm vis},$$
  
$$V_{\phi}(R,z) = V_K + \left(\frac{1}{St'^2 + 1} \frac{\rho_{\rm d}}{\rho_{\rm g} + \rho_{\rm d}} - 1\right) \eta V_K + \frac{St'}{St'^2 + 1} \frac{\rho_{\rm g}}{\rho_{\rm g} + \rho_{\rm d}} V_{\rm vis},$$

First term: Gas-dust friction, Second term: Gas viscosity  $\Omega_{\rm g} = \Omega_{\rm K}\sqrt{1-2\eta}, \qquad St' = \rho_{\rm g}/(\rho_{\rm g}+\rho_{\rm d})St$   $V_{\rm vis} = -\frac{3\nu}{R} \left[ p + \frac{2q}{3} + 2\left(\frac{z}{h_{\rm g}}\right)^2 \left(\frac{5q+9}{6}\right) \right] \qquad \eta(R,z) = -\frac{1}{2} \left(\frac{h_{\rm g}}{R}\right)^2 \left(p + q + \frac{q}{2}\frac{z^2}{h_{\rm g}^2}\right) \cdot \qquad St = t_{\rm stop}\Omega_{\rm K}$ 

Dust velocity

$$v_{\rm R}(R,z) = -\frac{2St'}{St'^2 + 1} \frac{\rho_{\rm g}}{\rho_{\rm g} + \rho_{\rm d}} \eta V_K + \frac{1}{St'^2 + 1} \frac{\rho_{\rm g}}{\rho_{\rm g} + \rho_{\rm d}} V_{\rm vis},$$
  
$$v_{\phi}(R,z) = V_K + \frac{1}{St'^2 + 1} \frac{\rho_{\rm g}}{\rho_{\rm g} + \rho_{\rm d}} \eta V_K - \frac{St'}{St'^2 + 1} \frac{\rho_{\rm g}}{\rho_{\rm g} + \rho_{\rm d}} V_{\rm vis},$$

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### Gas radial velocities in steady state

V<sub>vis</sub>: Viscous drift velocity, V<sub>K</sub>: Keplerian rotation velocity

Dipierro+17, Kanagawa+17



- Inward drift due to gas viscosity slows down due to dust feedback
- When the gas-dust friction term is larger than the viscous term, the gas can move outward against viscous diffusion.

### Drift velocities of gas

Solid: Vertical averaged velocity (3D), Dotted: 2D disk (ignore vertical structure)



# When can the dust feedback overcome viscosity ?

Above the line, the gas drift velocity is positive (gas move outward)



## Hydrodynamic simulations

- 2D (R, $\phi$ ) simulations by FARGO extended into gas-dust fluid.
  - Ignore the vertical structure of the disk (dust setting, etc)
- Commutating domain (6AU < R < 100AU) divided by radially 512 and azimuthally 128 meshes.
- Initial gas density  $\Sigma_g = 540 \text{ g/cm}^2 (\text{R}/1\text{AU})^{-1}$ .
- Initial dust/gas ratio,  $\Sigma_d / \Sigma_g = 0.01$ .
- Gas aspect ratio: h/R = 0.028 (R/1AU)<sup>1/4</sup>
- Constant size of dust grains (3cm)
  - Initially St=0.1 at 10AU
  - no growth, no fragmentation, no planetesimal formation
- Assuming that the disk size is large enough (dust and gas are supplied from the outside during the simulation)

### Effect on viscous disk evolution

Kanagawa+2017 ApJ

 $\times 10^{-5}$  $10^{3}$  $t = 3.2 \times 10^4 \text{ yr}$  $= 3.2 \times 10^4 \text{ yr}$ 6 Gas radial velocity  $t = 9.6 \times 10^4 \text{ yr}$ ensitv  $t = 9.6 \times 10^4 \text{ yr}$ Gas moves outward....  $t = 1.6 \times 10^5 \text{ yr}$  $t = 1.6 \times 10^5 \text{ yr}$ ..... 5 $10^{2}$  $V_R$  (AU/yr)  $\Sigma_g$  (g/cm<sup>2</sup>) surface 3  $10^{1}$ 2 Gas Gas density decreases because the gas flows outward  $10^{0}$  $10^{1}$  $10^{2}$  $10^{1}$  $10^{2}$ R (AU) R(AU)

 $\alpha = 10^{-3}$ ,  $\Sigma_d / \Sigma_g = 0.01$ , dust size = 3cm (initially St =0.1 at 10 AU)

### Implication for planet formation



But, ... we did not consider

- Size distribution of dust (Dipierro+18)
- 3D simulations (Gonzalez+15,17)
- Dust growth/fragmentation and planetesimal formation (Drążkowska+16, Drążkowska&Alibert 17)

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**Dust-to-Gas mass ratio** 

### Dust feedback can change a shape of a planet-induced gap When a large enough protoplanet forms, it would be a giant planet which forms a gap.

Dust size: 3 cm (constant),  $\alpha = 4 \times 10^{-3}$ , h/R=0.05, M<sub>p</sub>/M<sub>\*</sub>=10<sup>-3</sup>



Dust grains are highly accumulated at the outer edge  $\rightarrow$ Dust feedback must be effective.

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# Effect of dust feedback on migration of a gap-opening planet



Since the gas density at the outer disk decreases due to the dust feedback, torque exerted from the outer disk decreases.

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# Effect of dust feedback on migration of a gap-opening planet



Because of weaker negative torque from the outer disk, the planet can migrate outward.

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### Summary

- Dust feedback can help to make planets and save them from fast inward migration.
- The gas surface density within 10 AU can decrease due to the dust feedback.
  - Since the dust-to-gas mass ratio decreases, planetesimals and protoplanets could easy to be formed.
  - Since Type I migration is inefficient, the planets which are formed in this region can avoid too fast inward migration.
- Dust feedback reduce the torque exerted from the outer disk
  - The gap-opening planet can migrate outward when the planet is massive (Jupiter size) and viscosity is relatively low.
  - The gap-opening planet can survive for a long time.